# Stability of Simultaneously Triangularizable Switched Linear Systems on Time Scales 

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#### Abstract

Switched linear systems are a class of dynamical systems in which linear, time-invariant (LTI) dynamics closely model system behavior for a period of time, after which the system parameters suddenly (discontinuously) change, or "switch," and the system continues to evolve under new LTI parameters until another switching instant. Switched linear systems are often found wherever a dynamical system is coupled with supervisory control logic that can abruptly change the operating mode of the system, e.g. the transmission of a vehicle or the regulatory dynamics of a biological cell ${ }^{3}$. Until now, it has always been assumed that the underlying dynamics evolve naturally on $\mathbb{R}$ or $h \mathbb{Z}(h>0)$. However, many systems constructed with complex layers of computers and communication networks do not fit the standard choices for time domain, so recently, researchers ${ }^{1,2}$ have been investigating the use of dynamic equations on time scales to enable system modeling on unusual time domains $\mathbb{T}$, where $\mathbb{T}$ is any closed subset of $\mathbb{R}$. In this paper, we examine the stability of switched systems on discrete time scales, with arbitrary switching instances. We show that there exists a temporal region of stability for switched systems with simultaneously triangularizable system matrices, and that this region is identical to the region for simultaneously diagonalizable system matrices. It is henceforth assumed that the reader has a working knowledge of the field of time scale mathematics.


## Introduction

Consider the set of matrices $A_{1}, A_{2}, \ldots, A_{m} \in \mathbb{R}^{n \times n}$ with switching signal $s: \mathbb{T} \rightarrow 1,2, \ldots, m$, where

$$
\begin{equation*}
x^{\Delta}(t)=A_{s(t)} x(t) \quad \text { where } t \in \mathbb{T}, t \geq 0 \text { and } x(0)=x_{0} \tag{1}
\end{equation*}
$$

We make the following assumptions about this system and the underlying time scale
A1 The switching signal $s$ is arbitrary over $\mathbb{T}$.

A2 The eigenvalues of all of the $A_{i}$ are strictly within the Hilger circle for all $t \in \mathbb{T}$. (This means each $A_{i}$ is "stable," or $x^{\Delta}(t)=A_{i} x(t)$ has $\|x(t)\|<\infty$ for all $t \geq 0$.)

A3 All of the $A_{i}$ 's are regressive. (Meaning that $\left(I+\mu(t) A_{i}\right)^{-1}$ exists $\forall t \in \mathbb{T}$.)
A4 All of the $A_{i}$ 's commute, i.e. $A_{i} A_{j}=A_{j} A_{i} \forall i \neq j$.
A5 $\mathbb{T}$ has the following properties: (i) $0 \in \mathbb{T}$, (ii) $\mathbb{T}$ is unbounded above, (iii) $\mathbb{T}$ has graininess $0<\mu_{\min } \leq \mu(t) \leq \mu_{\max }$ for all $t \in \mathbb{T}$.

Note that all quantities can be assumed to be time-varying except for the $A_{i}$, unless otherwise indicated. Without loss of generality, we restrict the following analysis to $m=2$.

To investigate stability of (1), we define a Lyapunov candidate

$$
\begin{equation*}
V=x^{*} P x \tag{2}
\end{equation*}
$$

where $P=P^{*}>0$ and $*$ denotes conjugate transpose. To ensure stability, we need

$$
\begin{align*}
& V^{\Delta}<0 \\
\Rightarrow \quad & A_{i}^{*} P+P A_{i}+\mu A_{i}^{*} P A_{i}+\left(I+\mu A_{i}^{*}\right) P^{\Delta}\left(I+\mu A_{i}\right)<0 . \tag{3}
\end{align*}
$$

We now set

$$
\begin{equation*}
A_{i}^{*} P+P A_{i}+\mu A_{i}^{*} P A_{i}=-M_{i} \tag{4}
\end{equation*}
$$

where $M_{i}=M_{i}^{*}>0$ and solve for $P$. It can be shown ${ }^{4}$ that $P$ solves (4) for all $i$ if

$$
\begin{equation*}
P(t)=\int_{\mathbb{S}} \Phi_{A_{i}}(s, 0)^{*} M_{i}(t) \Phi_{A_{i}}(s, 0) \Delta s \tag{5}
\end{equation*}
$$

and

$$
\begin{align*}
& M_{1}(t)=\int_{\mathbb{S}} \Phi_{A_{2}}(s, 0)^{*} Q(t) \Phi_{A_{2}}(s, 0) \Delta s  \tag{6a}\\
& M_{2}(t)=\int_{\mathbb{S}} \Phi_{A_{1}}(s, 0)^{*} Q(t) \Phi_{A_{1}}(s, 0) \Delta s \tag{6b}
\end{align*}
$$

where $\mathbb{S}=\mu(t) \mathbb{N}_{0}, Q=Q^{*}>0$ is an arbitrary "seed" matrix, and $\Phi_{A_{i}}(s, 0)$ is the transition matrix that solves $y^{\Delta}(s)=A_{i} y(s)$ with $s \in \mathbb{S}$ and an initial condition $y(0)=y_{0}$. Because $\mathbb{S}$ is a constant-graininess time scale,

$$
\begin{equation*}
\Phi_{A_{i}}(s, 0)=e_{A_{i}}(s, 0)=\left(I+\mu(t) A_{i}\right)^{\frac{s}{\mu(t)}} . \tag{7}
\end{equation*}
$$

Substituting (4) into (3) yields

$$
\begin{equation*}
\left(I+\mu A_{i}\right)^{*} P^{\Delta}\left(I+\mu A_{i}\right)-M_{i}<0 \quad \text { for } t \in \mathbb{T} \tag{8}
\end{equation*}
$$

where $P^{\Delta}=\frac{P^{\sigma}-P}{\mu}$. Note that the only terms in (8) which depend on $t$ are $\mu$ and $\mu^{\sigma}$.
We can now pose the following question: given a time $t \in \mathbb{T}$, what is the region $\mathcal{R} \in \mathbb{R}^{2}$ such that $\left\{\mu(t), \mu^{\sigma}(t)\right\} \in \mathcal{R}$ satisfies (8) for all $i$ ?

Results have been derived ${ }^{2}$ for the diagonal and simultaneously diagonalizable cases, and we make note of them here. For simplicity, let

$$
\begin{align*}
K_{i} & :=\left[2 \operatorname{Re} J_{i}+\mu J_{i}^{*} J_{i}\right]  \tag{9a}\\
K_{i}^{\sigma} & :=\left[2 \operatorname{Re} J_{i}+\mu^{\sigma} J_{i}^{*} J_{i}\right] \tag{9b}
\end{align*}
$$

Then $\mathcal{R}$ is defined as the set of all $\left\{\mu(t), \mu^{\sigma}(t)\right\}$ such that

$$
\begin{equation*}
K_{1}^{\sigma} K_{2}^{\sigma} K_{1}^{-1} K_{2}^{-1}>\left(I+\mu J_{i}\right)^{*}\left(I+\mu J_{i}\right) \quad \text { for } i=1,2 \tag{10}
\end{equation*}
$$

with $0<\mu_{\min } \leq \mu(t) \leq \mu_{\max }$ for all $t \in \mathbb{T}$ and $A_{i}=S^{-1} J_{i} S$, where $J_{i}$ is diagonal. Note that $K_{i}^{-1}$ always exists because $K_{i}$ is diagonal and has non-zero eigenvalues as a result of A2.

Our objective is to investigate the simultaneously "triangularizable" case.

## Jordan Epsilon Form

We next explore the case where the $A_{i}$ 's are no longer simply diagonalizable. Thus, we need a more general form of the similarity transformation, $A=S^{-1} J S$. Therefore, let $A=S_{\epsilon}^{-1} J_{\epsilon} S_{\epsilon}$ with

$$
J_{\epsilon}=\left[\begin{array}{cccc}
\lambda & \epsilon & & \\
& \lambda & \epsilon & \\
& & \lambda & \epsilon \\
& & & \ddots
\end{array}\right]
$$

where $\epsilon$ is arbitrary and $0<\epsilon<1 . J_{\epsilon}$ comes from the generalized eigenvector equations

$$
\begin{aligned}
& A \bar{x}_{1}=\lambda \bar{x}_{1}, A \bar{x}_{2}=\lambda \bar{x}_{2}+\epsilon \bar{x}_{1}, A \bar{x}_{3}=\lambda \bar{x}_{3}+\epsilon \bar{x}_{2}, \ldots \\
\Rightarrow & (A-\lambda I) \bar{x}_{1}=0, \frac{(A-\lambda I)}{\epsilon} \bar{x}_{2}=\bar{x}_{1}, \frac{(A-\lambda I)}{\epsilon} \bar{x}_{3}=\bar{x}_{2}, \ldots \\
\Rightarrow & (A-\lambda I) \bar{x}_{1}=0, \frac{(A-\lambda I)^{2}}{\epsilon} \bar{x}_{2}=0, \frac{(A-\lambda I)^{3}}{\epsilon^{2}} \bar{x}_{3}=0, \ldots
\end{aligned}
$$

Thus, let $S_{\epsilon}=\left[\begin{array}{lll}\bar{x}_{1} & \epsilon \bar{x}_{2} \epsilon^{2} \bar{x}_{3} \ldots\end{array}\right]$ where the $\bar{x}_{i}$ are the eigenvectors of $A$. This is a valid similarity transform (as multiplication by a scalar doesn't change the fact that the $\bar{x}_{i}$ are eigenvectors). We term $J_{\epsilon}$ the "Jordan-epsilon form." From here on, let $S=S_{\epsilon}$ and $J=J_{\epsilon}$.

It will be useful to look at $J$ as the sum of two matrices, $J=L+N$, where

$$
L=\left[\begin{array}{cccc}
\lambda & 0 & & \\
& \lambda & 0 & \\
& & \lambda & 0 \\
& & & \ddots
\end{array}\right], \quad N=\left[\begin{array}{cccc}
0 & \epsilon & & \\
& 0 & \epsilon & \\
& & 0 & \epsilon \\
& & & \ddots
\end{array}\right]
$$

$N$ is a nilpotent matrix with $\epsilon$ on the superdiagonal, and $L$ is a diagonal matrix of the eigenvalues. Note that $N$ being nilpotent of order $n$ means that $N^{n}=[0]$, where $N($ and $A, L$, and $J)$ is $n \times n$.

## Triangular Transition Matrix

Remembering that $s \in \mathbb{S}:=\mu(t) \mathbb{N}_{0}$, we may now rewrite (7) with $h=\mu(t)$ and $z=\frac{s}{h} \in \mathbb{N}_{0}$ to obtain

$$
\begin{equation*}
\Phi_{A}(s, 0)=e_{A}(s, 0)=(I+h A)^{\frac{s}{h}}=(I+h A)^{z} . \tag{11}
\end{equation*}
$$

Applying the Jordan-epsilon similarity transform to $A$ then gives

$$
\begin{align*}
\Phi_{A}(s, 0) & =(I+h A)^{z} \\
& =\left(I+h S^{-1} J S\right)^{z} \\
& =S^{-1}\left(S S^{-1}+h J\right)^{z} S \\
& =S^{-1}(I+h J)^{z} S \\
& =S^{-1} \Phi_{J}(s, 0) S . \tag{12}
\end{align*}
$$

Expanding $\Phi_{J}$ as a binomial series we get

$$
\begin{align*}
\Phi_{J}(s, 0) & =(I+h J)^{z} \\
& =(I+h(L+N))^{z} \\
& =((1+h \lambda) I+h N)^{z} \\
& =\sum_{k=0}^{z}\binom{z}{k}(1+h \lambda)^{z-k}(h N)^{k} . \tag{13}
\end{align*}
$$

If $z<n-1$, then this is the final form of the binomial series, and the series is finite. If $z \geq n-1$, then the series is truncated at $n-1$ because $N$ is nilpotent, and we get

$$
\begin{align*}
= & (1+h \lambda)^{z} I+\binom{z}{1}(1+h \lambda)^{z-1}(h N)+\ldots \\
& \quad+\binom{z}{n-1}(1+h \lambda)^{z-n+1}(h N)^{n-1}+[0]+[0]+\ldots \\
= & (I+h L)^{z}+\sum_{k=1}^{n-1}\binom{z}{k}(1+h \lambda)^{z-k}(h N)^{k} \tag{14}
\end{align*}
$$

which is still finite for any $z$. Now, let

$$
\begin{equation*}
\Phi_{J}(s, 0)=(I+h L)^{z}+E(s)=e_{L}(s, 0)+E(s) \tag{15}
\end{equation*}
$$

where

$$
\begin{equation*}
E(s):=\sum_{k=1}^{z}\binom{z}{k}(1+h \lambda)^{z-k}(h N)^{k} \quad\left(\text { recall } z=\frac{s}{h}\right) . \tag{16}
\end{equation*}
$$

In other words, $E$ is an error matrix that depends on $\epsilon$, is upper-triangular with zeros on the diagonal, and is at most of order $\epsilon^{n-1}$. Note that $\|E\| \rightarrow 0$ as $\epsilon \rightarrow 0$.

## Error Terms

In the following discussion, many such error terms will appear, so we list them here for the sake of clarity:

$$
\begin{align*}
E_{i}(s) & :=\sum_{k=1}^{z}\binom{z}{k}\left[1+\mu \lambda_{i}\right]^{z-k}[\mu N]^{k}  \tag{17a}\\
E 1_{i} & :=\int_{\mathbb{S}}\left[e_{\bar{L}_{i}}(s, 0) E_{i}(s)+E_{i}^{*}(s) e_{L_{i}}(s, 0)+E_{i}^{*}(s) E_{i}(s)\right] \Delta s  \tag{17b}\\
E 2 & :=-E 1_{1}\left[2 \operatorname{Re} L_{2}+\mu L_{2}^{*} L_{2}\right]^{-1}-E 1_{2}\left[2 \operatorname{Re} L_{1}+\mu L_{1}^{*} L_{1}\right]^{-1}+E 1_{1} E 1_{2}  \tag{17c}\\
E 3_{i} & :=\mu\left(I+\mu J_{i}\right)^{-*} E 1_{j}\left(I+\mu J_{i}\right)^{-1}-E 2^{\sigma}+E 2  \tag{17d}\\
E 4_{i} & :=\mu\left(1+\mu \lambda_{i}\right) N^{*}+\mu\left(1+\mu \bar{\lambda}_{i}\right) N+\mu^{2} N N^{*}  \tag{17e}\\
E 5_{i} & :=E 4_{i}\left[\left(I+\mu J_{i}\right)\left(I+\mu J_{i}^{*}\right)\right]^{-1}+K K_{1} E 3_{i} K_{2} \tag{17f}
\end{align*}
$$

where $j=2,1$. Note: $E 2$ is the same for $i=1,2$.

## Generalizing the Region of Stability

Substituting (6a) into (5) for $i=1$, we get

$$
\begin{align*}
P & =\int_{\mathbb{S}} \Phi_{A_{1}}(s, 0)^{*} M_{1} \Phi_{A_{1}}(s, 0) \Delta s \\
& =\int_{\mathbb{S}} \Phi_{A_{1}}(s, 0)^{*}\left(\int_{\mathbb{S}} \Phi_{A_{2}}(r, 0)^{*} Q \Phi_{A_{2}}(r, 0) \Delta r\right) \Phi_{A_{1}}(s, 0) \Delta s \\
& =\int_{\mathbb{S}} \int_{\mathbb{S}} \Phi_{A_{1}}(s, 0)^{*} \Phi_{A_{2}}(r, 0)^{*} Q \Phi_{A_{2}}(r, 0) \Phi_{A_{1}}(s, 0) \Delta r \Delta s . \tag{18}
\end{align*}
$$

Since $Q$ in (6) may be any arbitrary, positive definite matrix, we choose $Q=S^{*} S$. Substituting this and applying the Jordan-epsilon similarity transform $A_{i}=S^{-1} J_{i} S$ to (18) gives

$$
\begin{align*}
P & =\int_{\mathbb{S}} \int_{\mathbb{S}}\left(S^{*} \Phi_{J_{1}}(s, 0)^{*} \Phi_{J_{2}}(r, 0)^{*} S^{-T}\right) S^{*} S\left(S^{-1} \Phi_{J_{2}}(r, 0) \Phi_{J_{1}}(s, 0) S\right) \Delta r \Delta s \\
& =S^{*}\left[\int_{\mathbb{S}} \Phi_{J_{1}}(s, 0)^{*} \Phi_{J_{1}}(s, 0) \Delta s \int_{\mathbb{S}} \Phi_{J_{2}}(r, 0)^{*} \Phi_{J_{2}}(r, 0) \Delta r\right] S . \tag{19}
\end{align*}
$$

We now need to investigate the integrals in this equation.

## Transition Matrix Quadratic Integral

Lemma 1 (Quadratic Integral). Let $A_{i}=S^{-1} J_{i} S$, where $J_{i} \in \mathbb{C}^{n \times n}$ is a Jordan Epsilon Form matrix, and be stable on time scale $\mathbb{S}=h \mathbb{N}_{0}$. Then

$$
\begin{equation*}
\int_{\mathbb{S}} \Phi_{J_{i}}(s, 0)^{*} \Phi_{J_{i}}(s, 0) \Delta s=-\left[2 \operatorname{Re} L_{i}+h L_{i}^{*} L_{i}\right]^{-1}+E 1_{i} \tag{20}
\end{equation*}
$$

where $J_{i}=L_{i}+N$, with $L_{i}$ diagonal and $N$ nilpotent, and $E 1_{i}$ is defined as in (17b).
Proof. Using the definition of $\Phi_{J}$ from (15) gives

$$
\begin{equation*}
\int_{\mathbb{S}} \Phi_{J_{i}}(s, 0)^{*} \Phi_{J_{i}}(s, 0) \Delta s=\int_{\mathbb{S}}\left[e_{L_{i}}^{*}(s, 0)+E_{i}^{*}(s)\right]\left[e_{L_{i}}(s, 0)+E_{i}(s)\right] \Delta s \tag{21a}
\end{equation*}
$$

By assumption A2, the left hand side of (21a) converges. We can expand it to see

$$
\begin{align*}
& =\int_{\mathbb{S}}\left[e_{\bar{L}_{i}}(s, 0) e_{L_{i}}(s, 0)+e_{\bar{L}_{i}}(s, 0) E_{i}(s)+E_{i}^{*}(s) e_{L_{i}}(s, 0)+E_{i}^{*}(s) E_{i}(s)\right] \Delta s \\
& =\int_{\mathbb{S}} e_{\bar{L}_{i}}(s, 0) e_{L_{i}}(s, 0) \Delta s+E 1_{i} \tag{21b}
\end{align*}
$$

where $E 1_{i}$ is defined as in (17b). $E 1_{i}$ must be finite because the integral on the left side of (21a) converges and the first term of (21b) is positive definite.

The first term of (21b) is an integral of a diagonal matrix. We can use the diagonal Quadratic Integral Lemma ${ }^{2}$ to say that

$$
\begin{equation*}
\int_{\mathbb{S}} e_{\bar{L}_{i}}(s, 0) e_{L_{i}}(s, 0) \Delta s=-\left[2 \operatorname{Re} L_{i}+h L_{i}^{*} L_{i}\right]^{-1} \tag{22}
\end{equation*}
$$

Substituting this into (21), we have

$$
\begin{equation*}
\int_{\mathbb{S}} \Phi_{J_{i}}(s, 0)^{*} \Phi_{J_{i}}(s, 0) \Delta s=-\left[2 \operatorname{Re} L_{i}+h L_{i}^{*} L_{i}\right]^{-1}+E 1_{i} \tag{23}
\end{equation*}
$$

which is the lemma statement.
The result of Lemma 1 closely parallels the result for the diagonalizable systems. In other words,

$$
\begin{equation*}
\lim _{\epsilon \rightarrow 0} \int_{\mathbb{S}} \Phi_{J_{i}}(s, 0)^{*} \Phi_{J_{i}}(s, 0) \Delta s=-\left[2 \operatorname{Re} L_{i}+h L_{i}^{*} L_{i}\right]^{-1} \tag{24}
\end{equation*}
$$

## The Region of Stability

Motivated by (9), we generalize $K_{i}$ by redefining it as

$$
\begin{align*}
K_{i} & :=\left[2 \operatorname{Re} L_{i}+\mu L_{i}^{*} L_{i}\right]  \tag{25a}\\
K_{i}^{\sigma} & :=\left[2 \operatorname{Re} L_{i}+\mu^{\sigma} L_{i}^{*} L_{i}\right] \tag{25b}
\end{align*}
$$

(recall that, in (9), the $J_{i}$ were diagonal, while now, the $L_{i}$ are the diagonal of $J_{i}$; also, $K_{i}^{-1}$ still exists because of the same reasoning as before).

Theorem 2 (Region of Stability). Given $A_{i}=S^{-1} J_{i} S$ where the $J_{i}$ are Jordan-epsilon form matrices with $J_{i}=L_{i}+N$ and the $A_{i}$ are stable on time scale $\mathbb{T}$, there exists a region $\mathcal{R} \in \mathbb{R}^{2}$ consisting of pairs $\left\{\mu^{\sigma}, \mu\right\}$ such that

$$
\begin{equation*}
K_{1}^{-1} K_{2}^{-1} K_{1}^{\sigma} K_{2}^{\sigma}>\left(I+\mu L_{i}^{*}\right)\left(I+\mu L_{i}\right) \quad \text { for } i=1,2 \tag{26}
\end{equation*}
$$

with $0<\mu_{\text {min }} \leq \mu(t) \leq \mu_{\text {max }}$ for all $t \in \mathbb{T}$.
Proof. Beginning with (19), we use Lemma 1 and (25) to say

$$
\begin{align*}
P & =S^{*}\left[\int_{\mathbb{S}} \Phi_{J_{1}}(s, 0)^{*} \Phi_{J_{1}}(s, 0) \Delta s \int_{\mathbb{S}} \Phi_{J_{2}}(r, 0)^{*} \Phi_{J_{2}}(r, 0) \Delta r\right] S \\
& =S^{*}\left[\left(-\left[2 \operatorname{Re} L_{1}+\mu L_{1}^{*} L_{1}\right]^{-1}+E 1_{1}\right)\left(-\left[2 \operatorname{Re} L_{2}+\mu L_{2}^{*} L_{2}\right]^{-1}+E 1_{2}\right)\right] S \\
& =S^{*}\left[\left(-K_{1}^{-1}+E 1_{1}\right)\left(-K_{2}^{-1}+E 1_{2}\right)\right] S \\
& =S^{*}\left[K_{1}^{-1} K_{2}^{-1}+E 2\right] S \tag{27}
\end{align*}
$$

where $E 2$ is defined as in (17c). Similarly

$$
\begin{equation*}
P^{\sigma}=S^{*}\left[K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}+E 2^{\sigma}\right] S \tag{28}
\end{equation*}
$$

Inserting $P$ and $P^{\sigma}$ from (27) and (28) into (8) for $i=1$ and eliminating $S$ gives

$$
\begin{gather*}
\frac{1}{\mu}\left(I+\mu J_{1}\right)^{*}\left[K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}+E 2^{\sigma}-K_{1}^{-1} K_{2}^{-1}-E 2\right]\left(I+\mu J_{1}\right)+K_{2}^{-1}-E 1_{2}<0  \tag{29}\\
\frac{1}{\mu}\left(I+\mu J_{1}\right)^{*}\left[K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}-K_{1}^{-1} K_{2}^{-1}+E 2^{\sigma}-E 2\right]\left(I+\mu J_{1}\right)<-K_{2}^{-1}+E 1_{2} \\
K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}-K_{1}^{-1} K_{2}^{-1}+E 2^{\sigma}-E 2<\mu\left(I+\mu J_{1}\right)^{-*}\left(-K_{2}^{-1}+E 1_{2}\right)\left(I+\mu J_{1}\right)^{-1} \\
K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}-K_{1}^{-1} K_{2}^{-1}<-\mu\left(I+\mu J_{1}\right)^{-*}\left(I+\mu J_{1}\right)^{-1} K_{2}^{-1}+E 3_{1} \tag{30}
\end{gather*}
$$

where $E 3$ is defined as in (17d). Continuing,

$$
\begin{align*}
K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}-K_{1}^{-1} K_{2}^{-1} & <-\mu\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)\right]^{-1} K_{2}^{-1}+E 3_{1} \\
K_{1} K_{2} K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}-I & <-\mu\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)\right]^{-1} K_{1}+K K_{1} E 3_{1} K_{2} \\
K_{1} K_{2} K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}} & <I-\mu\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)\right]^{-1} K_{1}+K_{1} E 3_{1} K_{2} \\
K_{1} K_{2} K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}} & <\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)-\mu K_{1}\right]\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)\right]^{-1}+K K_{1} E 3_{1} K_{2} \tag{31}
\end{align*}
$$

Multiplying $\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)$ out gives

$$
\begin{align*}
\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right) & =I+2 \mu \operatorname{Re} L_{1}+\mu^{2} L_{1} L_{1}^{*}+\mu\left(1+\mu \lambda_{1}\right) N^{*}+\mu\left(1+\mu \bar{\lambda}_{1}\right) N+\mu^{2} N N^{*} \\
& =I+\mu K_{1}+E 4_{1} \tag{32}
\end{align*}
$$

where we define $E 4$ as in (17e). Thus we have

$$
\begin{align*}
& K_{1} K_{2} K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}<\left[I+\mu K_{1}+E 4_{1}-\mu K_{1}\right]\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)\right]^{-1}+K_{1} E 3_{1} K_{2} \\
& K_{1} K_{2} K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}<\left[I+E 4_{1}\right]\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)\right]^{-1}+K_{1} E 3_{1} K_{2} \\
& K_{1} K_{2} K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}<\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)\right]^{-1}+E 4_{1}\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)\right]^{-1}+K_{1} E 3_{1} K_{2} \tag{33}
\end{align*}
$$

Let $E 5$ be defined as in (17f). Then

$$
\begin{align*}
& K_{1} K_{2} K_{1}^{\sigma^{-1}} K_{2}^{\sigma^{-1}}<\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)\right]^{-1}+E 5_{1} \\
& K_{1}^{-1} K_{2}^{-1} K_{1}^{\sigma} K_{2}^{\sigma}>>\left(\left[\left(I+\mu J_{1}\right)\left(I+\mu J_{1}^{*}\right)\right]^{-1}+E 5_{1}\right)^{-1} \\
& K_{1}^{-1} K_{2}^{-1} K_{1}^{\sigma} K_{2}^{\sigma}>>\left(\left[I+\mu K_{1}+E 4_{1}\right]^{-1}+E 5_{1}\right)^{-1} \\
& K_{1}^{-1} K_{2}^{-1} K_{1}^{\sigma} K_{2}^{\sigma}>\left(\left[\left(I+\mu L_{1}^{*}\right)\left(I+\mu L_{1}\right)+E 4_{1}\right]^{-1}+E 5_{1}\right)^{-1} \tag{34}
\end{align*}
$$

$i=2$ follows similarly.
We define a region $\mathcal{R}_{\epsilon}$ given by (34) for $i=1,2$. Then we note that

$$
\begin{equation*}
\lim _{\epsilon \rightarrow 0} \mathcal{R}_{\epsilon}=\mathcal{R} \tag{35}
\end{equation*}
$$

because

$$
\begin{equation*}
\lim _{\epsilon \rightarrow 0}\left(\left[\left(I+\mu L_{i}^{*}\right)\left(I+\mu L_{i}\right)+E 4_{i}\right]^{-1}+E 5_{i}\right)^{-1}=\left(I+\mu L_{i}^{*}\right)\left(I+\mu L_{i}\right) \tag{36}
\end{equation*}
$$

Applying this to (34) yields the set of equations in Theorem 2:

$$
\begin{equation*}
K_{1}^{-1} K_{2}^{-1} K_{1}^{\sigma} K_{2}^{\sigma}>\left(I+\mu L_{i}^{*}\right)\left(I+\mu L_{i}\right) \quad \text { for } i=1,2 \tag{37}
\end{equation*}
$$

Thus, to guarantee stability of the switched linear system (1), it is sufficient to show that $\left\{\mu(t), \mu^{\sigma}(t)\right\} \in$ $\mathcal{R}$, regardless of whether the system is simultaneously diagonalizable or only simultaneously triangularizable, which is the theorem statement.

This is a favorable result, because $\mathcal{R}$ is significantly easier to compute than $\mathcal{R}_{\epsilon}$.

## Examples

## Example 1

Let

$$
A_{1}=\left[\begin{array}{cc}
-1.2 & 0.2 \\
0 & -1.2
\end{array}\right], \quad A_{2}=\left[\begin{array}{cc}
-0.6 & 0.2 \\
0 & -0.6
\end{array}\right]
$$

which are already in irreducible, Jordan-epsilon form. Solving equation (37) for all $\left\{\mu(t), \mu^{\sigma}(t)\right\}$ pairs yields the region $\mathcal{R}$ in Figure 1, where it can be seen that the region is under the minimum of the four boundary curves (two curves exist outside the plot axes limits).


Figure 1: The temporal region of stability for Example 1. The two curves are the solution of (37) as if it were an equality, and the red area is under the minimum of the two curves. The upper limits on the axes are $\mu_{\max }$, which is the maximum graininess allowed by assumption A2.

## Example 2

Let

$$
A_{1}=\left[\begin{array}{cc}
-0.78 & 0.8 \\
0 & -0.78
\end{array}\right], \quad A_{2}=\left[\begin{array}{cc}
-0.1 & 0.8 \\
0 & -0.1
\end{array}\right] .
$$

Similarly to Example 1, solving (37) yields Figure 2.


Figure 2: The temporal region of stability for Example 2. The two curves are the solution of (37) as if it were an equality, and the red area is under the minimum of the two curves. The upper limits on the axes are $\mu_{\max }$, which is the maximum graininess allowed by assumption A2.

## References

1. DaCunha, J. "Stability for time varying linear dynamic systems on time scales," J. Comput. Appl. Math., Vol. 176, Issue 2, April 2005. 381-410.
2. Davis, J., Miller, J., Gravagne, I., Marks, R., Ramos, A. "Stabilty of Switched Linear Systems on Non-Uniform Time Domains," submitted to IEEE Trans. on Systems, Man, and Cybernetics - Part B. 2008.
3. Liberzon, D. Switching in Systems and Control. 2003, Birkhäuser, Boston.
4. Ramos, Alice. Stability of Hybrid Dynamical Systems: Analysis and Design. Ph.D. Thesis. Baylor University, Aug 2009.

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